

Short Communications

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The crystal structure of CsIOF₄.† By R. R. RYAN AND L. B. ASPREY, *Los Alamos Scientific Laboratory, University of California, Los Alamos, New Mexico 87544, U.S.A.*

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The colorless crystals of CsIOF₄ are orthorhombic with space group *Pmcn*, $Z=4$. Room-temperature cell constants are $a=6.024$, $b=15.488$, and $c=6.214$ Å. Data were collected on an automated four-circle diffractometer. The structure contains IOF₄⁻ anions (idealized symmetry C_{4v}) with four fluorine atoms in the equatorial plane and the oxygen atom on the axial site. The cesium ion has a coordination number of 10.

Recrystallization of stoichiometric amounts of CsI and IF₅ from acetonitrile produced colorless single crystals with the proper unit-cell volume to contain four molecules of CsIF₆. The Raman spectrum of the substance showed, however, an intense peak at 888 wave numbers which is expected for an I=O stretching frequency. Since the only ion which contains an I=O bond and has the right molecular volume for IF₆⁻ is IOF₄⁻, and since neither chemical nor structural information existed in the literature for this ion, an X-ray investigation was undertaken.

Several crystals were placed in glass capillaries in an inert atmosphere and examined by precession techniques. The crystal faces were of the prism {011} and of the pinacoid {100}. The photographs taken with Mo $K\alpha$ radiation showed orthorhombic symmetry with the systematic absences: $h0l$, l odd; $hk0$, $h+k$ odd. The systematic absences are consistent with the space groups *Pmcn* and *P2₁cn*. In all cases, the a^* axis was approximately along the axis of the glass capillary. ω scans of several reflections on an automated Picker diffractometer showed that all the crystals were fractured about the a^* axis so that the splitting of the peaks was most severe for large k and/or l values. The best crystal showed three peaks of approx-

imately equal intensity for which the maxima of the 'satellite' peaks were separated from the central peak by about 0.1° in ω for the worst reflections investigated. This crystal was used for determining cell constants and for data collection.

Cell dimensions and their estimated standard errors obtained by least-squares refinement on the positions of 12 high-order reflections as determined on the diffractometer using Mo $K\alpha_1$ radiation ($\lambda=0.70930$ Å) are $a=6.024$ (2), $b=15.488$ (6), and $c=6.214$ (3) Å.

Intensities were collected using Mo $K\alpha$ radiation, a single-crystal graphite monochromator (002 plane), and a take-off angle of 5° . The count was taken using a θ - 2θ scan over a 2θ range of 2° in 0.05° steps for two sec at each step; stationary-counter stationary-crystal background counts of 20 sec were taken at each end of the scan. Intensities were measured for all hkl planes in the positive octant and for their Friedel related reflections, resulting in 2315 measurements. The redundant (assuming *Pmcn*) reflections were averaged, resulting in 1003 reflections observed according to the criterion $I \geq 3\sigma(I)$ where $\sigma^2(I) = f^2(t)(T+B) + \sigma_b^2(T-B)^2$, T being the total count, B the estimated background, and $f(t)$ a time-dependent correction factor determined as described below. The quantity σ_b^2 was taken to be 2.5×10^{-4} , a number which we feel to be a reasonable estimate of the normalized variance of a re-

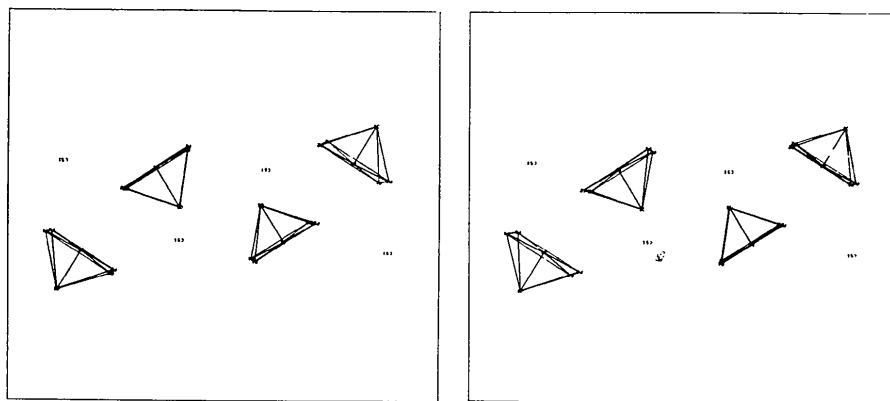


Fig. 1. Stereo view of the contents of the unit cell looking down the x axis with the y axis to the right. The four fluorine atoms near $y=\frac{1}{2}$ (and their symmetry equivalents) are F(1).

Table 1. Structure factors
Column headings are l, F0 and Fc

Table with multiple columns of numerical data representing structure factors for various reflections. The columns are organized into groups labeled with 'h' and 'k' values (e.g., h=0, k=0, 1, 2, 3, 4, 5, 6, 7, 8, 9, 10, 11, 12, 13, 14, 15, 16, 17, 18, 19, 20, 21, 22, 23, 24, 25, 26, 27, 28, 29, 30, 31, 32, 33, 34, 35, 36, 37, 38, 39, 40, 41, 42, 43, 44, 45, 46, 47, 48, 49, 50, 51, 52, 53, 54, 55, 56, 57, 58, 59, 60, 61, 62, 63, 64, 65, 66, 67, 68, 69, 70, 71, 72, 73, 74, 75, 76, 77, 78, 79, 80, 81, 82, 83, 84, 85, 86, 87, 88, 89, 90, 91, 92, 93, 94, 95, 96, 97, 98, 99, 100).

flexion due to random variations other than counting statistics. The R index, estimated by comparison of equivalent reflections, was 0.06.

The intensity of a standard reflection, measured after every 50 reflections was found to decrease by approximately 10% during the course of data collection. Multiplicative correction factors [f(t)] for the data were determined by fitting a polynomial to the intensity measurements of the standard reflection. The order of the polynomial was determined from significance tests based on the sum-of-the-squares of the residuals (Hamilton, 1965). Corrections were made for Lorentz and polarization factors. The graphite crystal in the monochromator was assumed to be of perfect mosaicity giving the following form for the polarization factor: (cos^2 2alpha + cos^2 2theta)/(1 + cos^2 2alpha), where alpha is the angle the beam makes with the monochromator. The absorption corrections were made with the Busing & Levy (1957) method using Burnham's (1962) program as modified by Larson, Cromer & Roof (1964). The linear absorption correction for Mo Kalpha radiation is 112 cm^-1, and the calculated transmission varied from 0.1 to 0.3.

Structure factors (Table 1) were calculated using the scattering factors of Doyle & Turner (1968) for the neutral atoms cesium, iodine, oxygen, and fluorine, and the dispersion terms of Cromer & Liberman (1970), i.e., Cs: Delta f' = -0.664, Delta f'' = 2.119 and I: Delta f' = -0.726 and Delta f'' = 1.812.

The Patterson function could be interpreted in terms of two heavy atoms in the special position 4(c) in Pmcn, one with the coordinates x approx 0.25, y approx 0.9, z approx 0.25, and one with the coordinates x approx 0.25, y approx 0.16, and z approx 0.25. Refinement of this structure, using iodine scattering amplitudes for both atoms, converged to an R index [R = sum(|Fo| - |Fc|) / sum |Fo|] of 0.15. A difference Fourier map showed three unique peaks. Two of these were about 2.0 A from a heavy atom and consequently were assigned as fluorine atoms. The remaining peak was assigned as an oxygen atom. The heavy atom that was surrounded by four fluorine atoms and one oxygen atom was assigned as iodine. The remaining heavy atom was assigned as cesium.

All refinements were carried out using Larson's (1971) full-matrix least-squares program. The function minimized was sum wi (|Fo| - |Fc|)^2, where wi is the weight defined as 1/sigma^2(Fo).

The structure refined to a weighted Rw index, Rw = sum wi^2 |Fo| - |Fc| / sum wi^2 |Fo| of 0.059 with all atoms anisotropic. The unweighted R index was 0.065. The Rw index with only the heavy atoms anisotropic was 0.070. The hypothesis that the light atoms vibrate isotropically can thus be rejected at the 0.05 confidence level. Parameter shifts in the last refinement were less than 0.001 sigma for all parameters. The standard deviation of an observation of unit weight is 4.7, indicating significant systematic errors in the data and/or model. The final difference Fourier map showed residuals near the heavy-atom position ranging in value from -2.0 to 2.3 e.A^-3, the largest residuals elsewhere were about 1.0 e.A^-3 with sigma(theta) = 0.4 e.A^-3. The final parameters are given in Table 2.

The structure of the IOF4- ion (see Fig. 1) has idealized symmetry C4v, with the oxygen atom occupying the axial position as is expected from electron-pair repulsion theory (Gillespie & Nyholm, 1957). This geometry has been established by microwave spectroscopy for the isoelectronic XeOF4 molecule (Martins & Wilson, 1968) with the parameters Xe-O = 1.703 (15), Xe-F = 1.900 (5) A and angle O-Xe-F = 91.8 (5) degrees. These parameters may be compared

Table 2. Fractional coordinates and thermal parameters ($\times 10^4$) for CsIOF₄

The thermal parameters are defined by the equation:

$$\exp [-(\beta_{11}h^2 + \beta_{22}k^2 + \beta_{33}l^2 + \beta_{12}hk + \beta_{13}hl + \beta_{23}kl)].$$

	<i>x</i>	<i>y</i>	<i>z</i>	β_{11}	β_{22}	β_{33}	β_{12}	β_{13}	β_{23}
Cs	2500	8968 (1)	2185 (2)	136 (2)	27.3 (4)	207 (3)	0	0	-2 (2)
I	2500	1609 (1)	2284 (1)	126 (2)	24.1 (2)	136 (2)	0	0	-7 (1)
O	2500	1030 (10)	9920 (20)	153 (23)	44 (6)	231 (32)	0	0	-107 (22)
F(1)	230 (10)	840 (10)	3450 (20)	298 (21)	42 (4)	384 (27)	-87 (16)	283 (43)	8 (2)
F(2)	180 (10)	2360 (10)	1020 (10)	162 (14)	52 (4)	232 (18)	54 (12)	-49 (29)	1 (2)

with the average distances and angles for IOF₄⁻ in the present study: I-F=1.965 (2), I-O=1.72 (1) Å and angle O-I-F=88.5 (2)°. Distances and angles for IOF₄ are shown in Table 3. The trend to shorter distances, as one moves to the right in the periodic table, has been established in the isoelectronic series MF₅E (see Table 4), although the magnitude of the change in the present case is certainly more than would be expected by such an analogy. In addition, it is tempting to argue that the larger O-M-F angle in the xenon compound (and the shorter distances) is at least partially due to the decrease in lone-pair to fluorine bond repulsions relative to the oxygen bond to fluorine bond repulsions as one moves to the right in the periodic table. It should be noted, however, that no such increase in angle can be established in the isoelectronic series shown in Table 4.

Table 3. Distances and angles

I—O	1.72 (1) Å	Cs—F(1)	(2)	3.18 (1)
I—F(1) (2)	1.95 (1)	Cs—F(1')	(2)	3.30 (1)
I—F(2) (2)	1.98 (1)	Cs—F(2)	(2)	3.29 (1)
O—F(1) (2)	2.60 (1)	Cs—F(2')	(2)	3.17 (1)
O—F(2) (2)	2.58 (1)	Cs—O	(2)	3.28 (1)
F(1)—F(2) (2)	2.80 (1)			
F(1)—F(1')	2.73 (1)	O—I—F(1) (2)		89.9 (4)°
F(2)—F(2')	2.79 (1)	O—I—F(2) (2)		88.1 (4)
		F(1)—I—F(2) (2)		90.7 (3)
		F(1)—I—F(1')		88.9 (5)
		F(2)—I—F(2')		88.6 (4)

The anion is coordinated by two fluorine atoms (at an I-F distance of 3.14 Å) which lie on a line perpendicular to the presumed iodine lone-pair direction and by three Cs⁺ ions which lie approximately above three of the triangular faces of the square pyramid. The intermolecular lone-pair

to fluorine atom interactions apparently lengthen the F(2)—F(2') distances, increase the F(2)—I—F(2') angles, and may also be primarily responsible for the lengthening of the I—F(2) distance with respect to I—F(1). The Cs⁺ ion is coordinated by eight fluorine atoms and two oxygen atoms at distances less than 3.5 Å.

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Table 4. Distances and angles in the isoelectronic series MF₅E

	SbF ₅ 2 ^{-*}	TeF ₅ ^{-*}	IF ₅ [†]	XeF ₅ [‡]
M—F (apical)	2.00 (9) Å	1.862 (4) Å	1.817 (10) Å	1.81 (8) Å
M—F (basal)	2.04 (9)	1.953 (3)	1.873 (5)	1.88 (8)
		1.952 (3)		
F _a —M—F _b	83 (2)°	79.4 (2)°	80.9 (2)°	79 (4)°
		78.3 (2)		80 (4)

* Mastin, Ryan & Asprey (1970).

† Jones, Burbank & Bartlett (1970).

‡ Bartlett, Einstein, Stewart & Trotter (1967).